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## Volumetric, viscometric and acoustic properties of binary mixtures of 2-propanol with *n*-alkanes  $(C_6, C_8, C_{10})$  at 298.15 and 308.15 K

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Densities ( $\rho$ ), viscosities ( $\eta$ ), and speeds of sound, ( $u$ ) of the binary mixtures of 2-propanol with *n*-alkanes (*n*-hexane, *n*-octane, and *n*-decane) were measured over the entire composition range at 298.15 and 308.15 K and at atmospheric pressure. Using the experimental values of density, viscosity and speed of sound, the excess molar volumes ( $V^{\text{E}}$ ), viscosity deviations ( $\Delta \eta$ ), deviations in speed of sound  $(\Delta u)$ , isentropic compressibility  $(\kappa_s)$ , deviations in isentropic compressibility ( $\Delta \kappa_s$ ), and excess Gibbs energies of activation of viscous flow ( $\Delta G^{*E}$ ) were calculated. These results were fitted to the Redlich–Kister type polynomial equation. The variations of these excess parameters with composition were discussed from the viewpoint of intermolecular interactions in these mixtures. The excess properties are found to be either positive or negative depending on the molecular interactions and the nature of liquid mixtures.

Keywords: densities; viscosities; excess molar volume; viscosity deviation; excess Gibbs energy of activation of viscous flow; speeds of sound; isentropic compressibility

### 1. Introduction

During the last few years the thermodynamic properties of binary mixtures of alcohol  $+$  alkane have been studied extensively  $[1–3]$ . The mixtures of alcohol with alkane are interesting for a number of scientific reasons, mostly connected to the formulation and testing of models to predict the properties of associating fluids, since the binary systems are composed of an associating liquid, the alcohol and a nonassociating one, the alkane [4]. Alcohol molecules associate mainly by building up linear chains. The extent of this association depends in a specific way on the concentration of the alcohol in the liquid state and on temperature [5]. On the other hand, alkanes are an important series of homologous, nonpolar organic solvents. They have often been used in the study of solute dynamics because their physico-chemical properties as a function of chain length are well-known [6]. The temperature dependence of the thermodynamic properties of these mixtures is also important, from the viewpoint of understanding the extent of selfassociation of alkanol molecules. The alkane  $+$  alkanol mixtures are also used as additives to petrol and in rectification processes for binary azeotropes [7].

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2-Propanol is widely used as coupling and dispersing agent in the chemical, pharmaceutical and household industries and as carrier and extraction solvent for natural products. The particular alkanes were chosen in view of their importance in hydrocarbon processing industries. The main purpose of our study is to characterize the molecular interactions in these systems and subsequently predict the effect of increasing chain length of alkanes and temperature on the molecular interactions.

The present article reports the experimental data for the densities,  $\rho$ , viscosities,  $\eta$ , and speeds of sound,  $u$  of 2-propanol with *n*-hexane, *n*-octane, and *n*-decane over the entire composition range at  $T = (298.15 \text{ and } 308.15) \text{ K}$  and at atmospheric pressure. Using the experimental values of  $\rho$ ,  $\eta$ , and u, the excess molar volumes  $(V^E)$ , viscosity deviations  $(\Delta \eta)$ , deviations in speed of sound  $(\Delta u)$ , isentropic compressibility  $(\kappa_s)$ , deviations in isentropic compressibility ( $\Delta \kappa_s$ ), and excess Gibb's energies of activation of viscous flow  $(\Delta G^{*E})$  were calculated. These results were fitted to the Redlich–Kister type polynomial equation [8]. An attempt is also made to rationalise the results by collecting the data on 2-propanol  $+n$ -alkane mixtures.

#### 2. Experimental

#### 2.1. Materials

2-Propanol was supplied by Qualigens Fine Chemicals Pvt. Ltd. (Mumbai, India). The analytical grade n-hexane, n-octane and n-decane were supplied by S.D. Fine Chemicals Ltd. The purities of these reagents are  $>99\%$  in mass. Before use, all the chemicals were degassed by vacuum pump. The experimentally measured values of  $\rho$ ,  $\eta$ ,  $u$ , and  $\kappa_s$  of all the pure components at 298.15 and 308.15 K with their corresponding literature values [9–23] are listed in Table 1. These comparisons indicated good consistency between these results and those presented in the literature.

#### 2.2. Apparatus and procedure

The binary mixtures were prepared by mass in special air-tight bottles. The weighing was done on an electronic balance with a precision of  $\pm 0.1$  mg. The densities were measured using a bicapillary pycnometer of about  $10 \text{ cm}^3$  volume and a capillary bore with an internal diameter of 1 mm. The pycnometer was calibrated by using double- distilled, deionized, and degassed water with 997.05 and 994.03 kg  $\text{m}^{-3}$  with densities at 298.15 and 308.15 K, respectively. The accuracy in density measurements was estimated as  $\pm 0.0003 \times 10^{3}$  kg m<sup>-3</sup>.

The kinematic viscosities,  $\nu$ , were measured by means of a suspended-level Ubbelohde viscometer. The time of flow was measured with a stopwatch capable of recording to 0.01 s. For each mixture, at least four flow time measurements were performed and the results were averaged. The viscometer was calibrated at the experimental temperatures using the relation  $\nu = \eta/\rho$  =  $(C_1t - C_2)/t$ , where  $\nu$  is the kinematic viscosity,  $\eta$  is the absolute viscosity,  $\rho$  is the density, and t is the flow time of calibration liquid.  $C_1$  and  $C_2$ are calibration constants of the viscometer. The estimated accuracy in  $\eta$  measurements is within  $\pm 0.003$  mPa s.

The speeds of sound were measured using ultrasonic interferometer (Model MX-3) provided by Mittal Enterprises, New Delhi. The instrument was calibrated by measuring the speeds of sound in standard liquids e.g., water, acetone, cyclohexane, ethanol, etc.

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Table 1. Comparison of experimental densities  $\rho$ , viscosities  $\eta$ , speeds of sound  $u$ , and isentropic compressibility  $\kappa_s$ , with literature data for pure Table 1. Comparison of experimental densities  $\rho$ , viscosities  $\eta$ , speeds of sound  $u$ , and isentropic compressibility  $\kappa_s$ , with literature data for pure liquids.



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The temperature of the test liquids and their binary mixtures, during the measurements, was maintained to an accuracy of  $\pm 0.01 \text{ K}$  by using an electronically controlled thermostatic water bath.

#### 3. Results and discussion

The values of density,  $\rho$ , viscosity,  $\eta$ , speed of sound, u, excess molar volumes,  $V^E$ , viscosity deviation,  $\Delta \eta$ , deviation in speed of sound,  $\Delta u$ , isentropic compressibility,  $\kappa_s$ , deviation in isentropic compressibility,  $\Delta \kappa_{\rm s}$ , and excess Gibbs energy of activation of viscous flow,  $\Delta G^{*\rm E}$ for the binary mixtures at 298.15 and 308.15 K are recorded in Tables 2 and 3, respectively.

The excess functions, which are a measure of deviation from ideal behaviour, are found to be highly sensitive to intermolecular interactions between component molecules of the mixture. The sign and magnitude of these functions depend on the strength of interaction between dissimilar molecules in the mixture. The excess functions such as  $V^E$ ,  $\Delta \eta$ ,  $\Delta u$ ,  $\Delta \kappa_s$ , and  $\Delta G^{*E}$  were calculated by using the following equations:

$$
V^{E} \text{ (m}^{3} \text{ mol}^{-1}) = \frac{(M_{1}x_{1} + M_{2}x_{2})}{\rho_{\text{m}}} - \left(\frac{M_{1}x_{1}}{\rho_{1}} + \frac{M_{2}x_{2}}{\rho_{2}}\right),\tag{1}
$$

$$
\Delta \eta \text{ (mPa s)} = \eta_{\text{m}} - (x_1 \eta_1 + x_2 \eta_2), \tag{2}
$$

$$
\Delta u \, (\text{m s}^{-1}) = u_{\text{m}} - (x_1 u_1 + x_2 u_2), \tag{3}
$$

$$
\Delta \kappa_{\rm s} = \kappa_{\rm s}^{\rm mix} - (\varphi_1 \kappa_{\rm s,1} + \varphi_2 \kappa_{\rm s,2}). \tag{4}
$$

Volume fraction,  $\varphi_i$  is defined by the relation:

$$
\varphi_i = \frac{x_i V_i}{\sum_{i=1}^2 x_i V_i},\tag{5}
$$

$$
\Delta G^{*E} = RT[\ln(\eta_m V) - x_1 \ln(\eta_1 V_1) - x_2 \ln(\eta_2 V_2)],
$$
\n(6)

where  $\rho_m$ ,  $\eta_m$ ,  $u_m$ ,  $\kappa_s^{mix}$ , and V denote the density, viscosity, speed of sound, isentropic compressibility, and molar volume of the mixtures, respectively.  $\rho_i$ ,  $\eta_i$ ,  $u_i$ ,  $\kappa_{s,i}$ ,  $V_i$ , and  $M_i$  $(i = 1 \text{ and } 2)$  denote, respectively, the density, viscosity, speed of sound, isentropic compressibility, molar volume, and molecular weight of i-th component.

Isentropic compressibility,  $\kappa_s$ , and molar volume, V, were calculated using the relation:

$$
\kappa_{\rm s} = (\rho u^2)^{-1},\tag{7}
$$

$$
V = \frac{M_1 x_1 + M_2 x_2}{\rho_m}.
$$
\n(8)

The values of  $V^E$ ,  $\Delta \eta$ ,  $\Delta u$ , and  $\Delta G^{*E}$  were fitted to Redlich–Kister type polynomial equation:

$$
Y = x_1 x_2 \sum_{i=1}^{k} A_i (x_1 - x_2)^{i-1},
$$
\n(9)

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Table 2. Experimental densities  $\rho$ , viscosities  $\eta$ , speeds of sound u, isentropic compressibility  $\kappa_s$ , excess molar volume  $V^E$ , viscosity deviation  $\Delta$ Table 2. Experimental densities  $\rho$ , viscosities  $\eta$ , speeds of sound  $u$ , isentropic compressibility  $\kappa_s$ , excess molar volume  $V^E$ , viscosity deviation  $\Delta \eta$ , deviation and deviation  $\Delta u$ , and deviation in isentr deviation in speeds of sound  $\Delta$ u, and deviation in isentropic compressibility  $\Delta$  $\kappa_s$  and excess Gibbs energies of activation of viscous flow,  $\Delta$  $G^*^E$  for the binary mixtures at 298.15 K.





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Table 3. Experimental densities  $\rho$ , viscosities  $\eta$ , speeds of sound u, isentropic compressibility  $\kappa_s$ , excess molar volume  $V^E$ , viscosity deviation  $\Delta$ лацоп ∆*η*,<br>G<sup>∗E</sup> for the deviation in speeds of sound  $\Delta$ u, and deviation in isentropic compressibility  $\Delta$  $k<sub>s</sub>$  and excess Gibb's energies of activation of viscous flow,  $\Delta$ binary mixtures at 308.15 K.





Table 4. Values of  $A_i$  parameters of Equations (9) and (11) along with standard deviations  $\sigma$  for excess properties of the binary mixtures at 298.15 K.

	Binary mixture Excess property	$A_0$	$A_1$	A <sub>2</sub>	$A_3$	$A_{4}$	$\sigma$
$2$ -Propanol $+$ $n$ -Hexane	$V^{\rm E} \times 10^6$ (m <sup>3</sup> mol <sup>-1</sup> )	1.732	$-1.110$	0.770	$-0.330$		0.003
	$\Delta \eta$ (mPas)	$-2.776$	$-1.671$	$-0.983$	$-1.059$	$-0.699$	0.003
	$\Delta u$ (m s <sup>-1</sup> )	$-145.4$	18.2	9.0			0.3
	$\Delta \kappa_s$ (TPa <sup>-1</sup> )	152.5	$-219.3$	118.2	$-34.3$		0.6
	$\Delta G^{*E}$ (J mol <sup>-1</sup> )	$-4445.6$	$-298.6$	116.9	$-1043.9$	$-462.4$	10.2
$2$ -Propanol $+$ $n$ -Octane	$V^{\rm E} \times 10^6$ (m <sup>3</sup> mol <sup>-1</sup> )	2.341	$-0.477$	0.612	$-0.729$		0.004
	$\Delta \eta$ (mPa s)	$-2.505$	$-1.493$	$-0.843$	$-1.073$	$-1.040$	0.003
	$\Delta u$ (m s <sup>-1</sup> )	$-59.8$	$-23.0$	39.3			0.3
	$\Delta \kappa_s$ (T Pa <sup>-1</sup> )	146.5	$-114.3$	11.9	83.4	$-67.9$	0.6
	$\Delta G^{*E}$ (J mol <sup>-1</sup> )	$-3645.8$	$-958.4$	$-152.8$	$-437.8$	$-1067.8$	3.5
$2$ -Propanol $+$ $n$ -Decane	$V^{\rm E} \times 10^6$ (m <sup>3</sup> mol <sup>-1</sup> )	2.503	$-0.295$	0.679	$-0.672$		0.004
	$\Delta \eta$ (mPa s)	$-2.117$	$-1.161$	$-0.551$	$-0.978$	$-1.198$	0.004
	$\Delta u$ (m s <sup>-1</sup> )	14.3	25.9	$-13.7$	14.8	32.2	0.3
	$\Delta \kappa_s$ (TPa <sup>-1</sup> )	103.3	$-55.1$	42.3	$-79.8$		0.7
	$\Delta G^{*E}$ (J mol <sup>-1</sup> )	$-2371.6$	$-959.5$	$-353.6$	$-401.9$	$-1094.6$	16.5

where Y is  $V^E$  or  $\Delta \eta$  or  $\Delta u$  or  $\Delta G^{*E}$ , k is the number of estimated parameters and  $A_i$ represents polynomial coefficients, which were obtained by fitting the equation to the experimental results with least-squares regression method. The corresponding standard deviations  $\sigma(Y)$  were calculated using the relation:

$$
\sigma(Y) = \left[\frac{\sum (Y_{\text{expt.}} - Y_{\text{calcd}})^2}{(n-k)}\right]^{1/2},\tag{10}
$$

where *n* is the no. of experimental data points. The  $\Delta \kappa$  values were correlated with composition data by relation

$$
\Delta \kappa_{\rm s} = \varphi_1 \varphi_2 \sum_{i=1}^k A_i (\varphi_1 - \varphi_2)^{i-1} . \tag{11}
$$

The standard deviations  $\sigma(\Delta \kappa_s)$  may be given as

$$
\sigma(\Delta \kappa_{\rm s}) = \left[ \frac{\sum \left( \Delta \kappa_{\rm s}(\exp) - \Delta \kappa_{\rm s}(\text{calcd}) \right)^2}{(n-k)} \right]^{1/2}.
$$
\n(12)

The parameters of Equations (9) and (11) and standard deviations of Equations (10) and (12) are listed in Tables 4 and 5. The variation of smoothed values of  $V^E$ ,  $\Delta \eta$ ,  $\Delta u$ , and  $\Delta G^{*E}$  with mole fraction  $x_1$  and that of  $\Delta \kappa_s$  with volume fraction  $\varphi_1$  of 2-propanol are shown graphically in Figures 1–5, respectively.

Figure 1(a) and (b) display the dependence of  $V^E$  on  $x_1$  for all the mixtures at 298.15 and 308.15 K, respectively. It is observed that, for all the mixtures, these values are positive, suggesting the weak dispersive-type intermolecular interactions. This can be

	Binary mixture Excess property	$A_0$	$A_1$	A <sub>2</sub>	$A_3$	$A_4$	$\sigma$
$2$ -Propanol $+$ $n$ -Hexane	$V^{\rm E} \times 10^6$ (m <sup>3</sup> mol <sup>-1</sup> )	2.274	$-0.893$	0.978			0.025
	$\Delta \eta$ (mPas)	$-1.887$	$-0.960$	$-0.372$	0.273		0.014
	$\Delta u$ (m s <sup>-1</sup> )	$-151.1$	12.5	43.9	49.4	$-49.4$	0.9
	$\Delta \kappa_s$ (T Pa <sup>-1</sup> )	159.3	$-269.2$	$-24.9$	$-65.0$	279.5	2.3
	$\Delta G^{*E}$ (J mol <sup>-1</sup> )	$-3202.8$	496.9	$-404.1$	1704.5		33.5
$2$ -Propanol $+$ $n$ -Octane	$V^{\rm E} \times 10^6$ (m <sup>3</sup> mol <sup>-1</sup> )	2.581	$-0.819$	1.084	$-1.003$		0.017
	$\Delta \eta$ (mPa s	$-1.799$	$-1.083$	$-0.731$	$-0.564$	$-0.396$	0.002
	$\Delta u$ (m s <sup>-1</sup> )	$-66.7$	$-15.7$	$-39.9$			0.7
	$\Delta \kappa_s$ (T Pa <sup>-1</sup> )	175.3	$-53.5$	48.5	$-97.9$	82.3	1.6
	$\Delta G^{*E}$ (J mol <sup>-1</sup> )	$-3259.9$	$-1010.5$	$-712.0$			6.1
$2$ -Propanol $+$ $n$ -Decane	$V^{\rm E} \times 10^6$ (m <sup>3</sup> mol <sup>-1</sup> )	2.940	$-0.730$	1.130	$-1.698$		0.020
	$\Delta \eta$ (mPas)	$-1.465$		$-0.778 -0.750$	$-0.858$		0.006
	$\Delta u$ (m s <sup>-1</sup> )	14.0	4.7	$-17.1$	34.5		0.7
	$\Delta \kappa_s$ (T Pa <sup>-1</sup> )	118.3	$-42.8$	$-46.8$	$-68.2$	163.9	1.8
	$\Delta G^{*E}$ (J mol <sup>-1</sup> )	$-1975.7$	$-692.9$	$-1253.1$	$-1337.2$	1318.1	20.4

Table 5. Values of  $A_i$  parameters of Equations (9) and (11) along with standard deviations  $\sigma$  for excess properties of the binary mixtures at 308.15 K.

explained by the fact that 2-propanol is hydrogen bonded self-associated liquid and the addition of alkanes, which are nonpolar and nonassociated, causes the rupture of hydrogen bonds in 2-propanol with resulting volume expansion and  $V<sup>E</sup>$  is positive. The values of  $V^E$  vary according to the sequence *n*-hexane  $\lt$  *n*-octane  $\lt$  *n*-decane. The nature of  $V^E$  versus  $x_1$  curves for all the mixtures at both the temperatures is almost identical but the  $V<sup>E</sup>$  values show an increase with rise in temperature.

The dependence of  $\Delta \eta$  on the mole fraction of 2-propanol at 298.15 and 308.15 K is displayed in Figure 2(a) and (b), respectively. We observe that  $\Delta \eta$  values are negative for all the binary mixtures, further supporting the presence of weak intermolecular interactions. With an increase in temperature, the absolute values of  $\Delta \eta$  decrease as observed in the present study. This can be explained by the fact that an increment of temperature diminishes the self association in the pure components (2-propanol) and also the hetero-association between unlike molecules, because of the increase in thermal energy. This leads to less negative values of  $\Delta \eta$  as temperature is raised.

Figure 3(a) and (b) show the variation of  $\Delta u$  with mole fraction of 2-propanol at 298.15 and 308.15 K, respectively. The  $\Delta u$  values show a negative trend for 2-propanol  $+n$ -hexane and 2-propanol  $+n$ -octane binary mixtures over the entire composition range, whereas for 2-propanol + *n*-decane binary mixture,  $\Delta u$  is negative at lower  $x_1$  values of 2-propanol but it shows a positive trend at higher  $x_1$  of 2-propanol. The negative  $\Delta u$  values show the importance of interstitial accommodation of alkane molecules between alcohol aggregates, whereas positive  $\Delta u$  values for 2-propanol + n-decane suggest that as the size of the alkane molecule increases, the molecules of the mixing component cannot accommodate easily.

The deviations in isentropic compressibility,  $\Delta \kappa_s$ , at 298.15 and 308.15 K are positive over the whole range of composition for all the binary mixtures formed by 2-propanol



Figure 1. (a) Excess molar volumes  $V^E$  at 298.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , nhexane;  $\bullet$ , *n*-octane;  $\blacktriangle$ , *n*-decane; (--) for Redlich–Kister Equation (9);  $\nabla$  represents  $V^E$  values for 2-propanol  $+n$ -octane; and  $\blacklozenge$  for 2-propanol  $+n$ -decane at 298.15 K given by Tojo *et al.* (b) Excess molar volumes  $V^E$  at 308.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , *n*-hexane;  $\bullet$ , *n*-octane;  $\blacktriangle$ , *n*-decane; and  $\left(\text{---}\right)$  for Redlich–Kister Equation (9).

with *n*-alkanes, as shown in Figure  $4(a)$  and (b), respectively. The positive deviations in  $\Delta \kappa$ <sub>s</sub> may be attributed to the presence of weak interactions between the component molecules in the mixture, thereby indicating the predominance of long range dispersive forces. The positive  $\Delta \kappa$  values also suggest that the mixture is more compressible than the



Figure 2. (a) Viscosity deviations  $\Delta \eta$  at 298.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , *n*-hexane;  $\bullet$ , *n*-octane;  $\blacktriangle$ , *n*-decane; (-) for Redlich–Kister Equation (9);  $\blacktriangledown$  represents  $\Delta \eta$  values for 2-propanol +n-octane; and  $\blacklozenge$  for 2-propanol +n-decane at 298.15 K given by Tojo et al. (b) Viscosity deviations  $\Delta \eta$  at 308.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , *n*-hexane;  $\bullet$ , *n*-octane;  $\blacktriangle$ , *n*-decane; and  $(-)$  for Redlich–Kister Equation (9).

corresponding ideal mixture. Therefore, in these binary mixtures, an expansion in free volume is considered to occur, making the mixtures more compressible than the ideal mixture. As the temperature increases, the absolute value of  $\Delta \kappa_s$  also increases in all the binary mixtures.



Figure 3. (a) Deviation in speed of sound  $\Delta u$  at 298.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , n-hexane;  $\bullet$ , n-octane;  $\blacktriangle$ , n-decane; and (-) for Redlich–Kister Equation (9). (b) Deviation in speed of sound  $\Delta u$  at 308.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , *n*-hexane;  $\bullet$ , *n*-octane;  $\blacktriangle$ , *n*-decane; and  $(-)$  for Redlich–Kister Equation (9).



Figure 4. (a) Deviation in isentropic compressibility  $\Delta \kappa_s$  at 298.15K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , n-hexane;  $\bullet$ , n-octane;  $\blacktriangle$ , n-decane; and (—) for Redlich–Kister Equation (11). (b) Deviation in isentropic compressibility  $\Delta \kappa_s$  at 308.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , *n*-hexane;  $\bullet$ , *n*-octane;  $\blacktriangle$ , *n*-decane; and (-) for Redlich–Kister Equation (11).



Figure 5. (a) Excess Gibbs energies of activation of viscous flow  $\Delta G^{*E}$  at 298.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , n-hexane;  $\bullet$ , n-octane;  $\blacktriangle$ , n-decane; and (—) for Redlich–Kister Equation (9). (b) Excess Gibbs energies of activation of viscous flow  $\Delta G^{*E}$  at 308.15 K for binary mixtures of 2-propanol (1) with  $\blacksquare$ , n-hexane;  $\bullet$ , n-octane;  $\blacktriangle$ , n-decane; and (-) for Redlich–Kister Equation (9).

Figure 5(a) and (b) show the variation of excess free energy of activation of viscous flow  $\Delta G^{*E}$  with the mole fraction  $x_1$  of 2-propanol. The  $\Delta G^{*E}$  values are negative over the entire composition range for all the binary mixtures but the absolute value of  $\Delta G^{*E}$  diminishes with increasing chain length of the alkanes. The order it follows is: *n*-hexane > *n*-octane > *n*-decane. The negative values of  $\Delta G^{*E}$  indicate the presence of

weak interactions between the mixing components. This behaviour is also supported by the positive values of  $V^E$ . As the temperature increases, the absolute values of  $\Delta G^{*E}$  decrease for all these binary mixtures.

In the literature, Tojo et al. [24] presented the values of  $V^E$  and  $\Delta \eta$  for the binary mixtures of 2-propanol with *n*-octane and *n*-decane at 298.15 K. Comparison of our  $V^E$ and  $\Delta \eta$  values at 298.15 K suggest a reasonably good agreement for these systems. The shapes of the curves also remain almost identical to our data. No experimental data of  $\rho$ ,  $\eta$ , and u on these mixtures are available in the literature at 308.15 K with which we can compare the present results. The  $V^E$  values for the binary mixture of 2-propanol + n-hexane at 298.15 K and 308.15 K have also been compared with  $V<sup>E</sup>$  results given by Modarress et al. [25]. The shapes of curves also remain almost identical to our data, but our values are lower than those given by Modarress et al.

#### 4. Conclusion

The interaction between 2-propanol and  $n$ -alkane is such that the basic network of intermolecular association in 2-propanol is disrupted. The behaviour of 2-propanol  $+$ n-alkane mixtures indicates the presence of long dispersive forces among the components. Positive deviations are observed in the case of  $V^E$  and  $\Delta \kappa_s$ , whereas  $\Delta \eta$  and  $\Delta G^{*E}$  show negative deviations for all binary mixtures. Both positive and negative deviations are observed in the case of  $\Delta u$ . The behaviour of all these excess parameters studied here support each other. On the whole, it can be concluded that the strength of bonding is expected to decrease with an increase in the chain length of alkane molecule and also with rise in temperature.

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